

More on Angular Momentum

$$\begin{aligned} L_{op}^2 &= -\hbar^2 \left[\frac{2}{2} + \cot^2 \theta \right] + \frac{1}{\sin^2 \theta} \frac{2}{2} \\ L_x &= i\hbar \left(\sin \theta \frac{\partial}{\partial \phi} + \cot \theta \cos \theta \frac{\partial}{\partial \theta} \right) \\ L_y &= -i\hbar \left(\cos \theta \frac{\partial}{\partial \phi} - \cot \theta \sin \theta \frac{\partial}{\partial \theta} \right) \\ L_z &= -i\hbar \frac{\partial}{\partial \phi} \end{aligned}$$

Eigenfunctions of L_{op}^2

First, $[L_z, L_{op}^2] = 0$ so eigs. of L_z are eigs of L_{op}^2 .

$$L_z Y_{\ell m}(\theta, \phi) = m\hbar Y_{\ell m}(\theta, \phi) \quad Y_{\ell m}(\theta, \phi) = \frac{e^{im\phi}}{\sqrt{2\pi}}$$

$$-i\hbar \frac{\partial}{\partial \phi} Y_{\ell m} = -i\hbar(im) Y_{\ell m}(\theta, \phi)$$

$$L_z Y_{\ell, m+2} = m\hbar Y_{\ell, m+2} \quad Y_{\ell, m+2}(\theta, \phi) = Y_{\ell, m}(\theta, \phi)$$

So eigenvalues of L_z are $m\hbar$, $m=0, \pm 1, \pm 2, \dots$

(recall plane rotor-- $\frac{L_z^2}{2\mu r_e^2}$)

$$Y_{\ell m}(\theta, \phi) = \frac{e^{im\phi}}{\sqrt{2\pi}} P_{\ell}^{-m}(\cos \theta) \quad \text{--Spherical Harmonic}$$

where the eigenvalue equation for $Y_{\ell m}$ is

$$L_{op}^2 Y_{\ell m} = Y_{\ell m} \left[-\hbar^2 \left(\frac{2}{2} + \cot^2 \theta \right) + \frac{\hbar^2 m^2}{\sin^2 \theta} - \frac{-\hbar^2 m^2}{\sin^2 \theta} \right] (\cos \theta) = 0$$

Solutions are known—Associated Legendre Polynomials

And $Y_{\ell m} = \hbar^2 \ell(\ell+1)$, $\ell = 0, 1, 2, \dots$ and $-\ell \leq m \leq \ell$.

Summary so far:

$$L_z Y_{\ell m}(\theta, \phi) = L_z \frac{1}{\sqrt{2}} Y_{\ell}^{-|m|}(\cos \theta) \frac{e^{im}}{\sqrt{2}} = \frac{1}{\sqrt{2}} Y_{\ell}^{-|m|}(\cos \theta) m \hbar \frac{e^{im}}{\sqrt{2}}$$

$$L_z Y_{\ell m} = m \hbar Y_{\ell m}(\theta, \phi)$$

$$L_{op}^2 Y_{\ell m} = \hbar^2 \ell(\ell + 1) Y_{\ell m}(\theta, \phi) \quad \ell = 0, 1, 2, \dots \quad \text{and} \quad -\ell \leq m \leq \ell$$

Orthonormality

$$\langle Y_{\ell m} | Y_{\ell' m'} \rangle = \delta_{\ell \ell'} \delta_{m m'}$$

$$\int_0^{2\pi} \int_{-1}^1 d(\cos \theta) Y_{\ell m}^* Y_{\ell' m'} = \delta_{\ell \ell'} \delta_{m m'}$$

(spherical polar coordinates: $dx dy dz = r^2 dr d(\cos \theta) d\phi$)

Integral over ϕ is straightforward.

$$Y_{\ell}^{-|m|}(\cos \theta) = \frac{2\ell + 1}{2} \frac{(l - |m|)!}{(l + |m|)!}^{1/2} Y_{\ell}^{|m|}(\cos \theta)$$

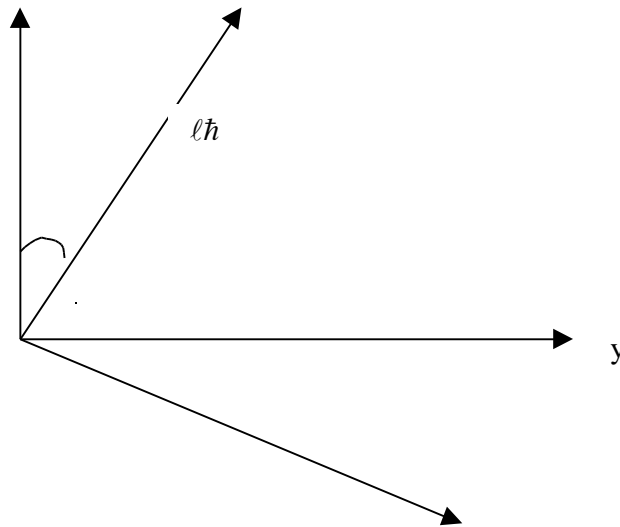
$$\text{and} \quad Y_{\ell}^{|m|}(x) = \frac{1}{2^{\ell} \ell!} (1 - x^2)^{|m|/2} \frac{d^{\ell + |m|}}{dx^{\ell + |m|}} (x^2 - 1)^{\ell} \quad \ell = 0, 1, 2, \dots$$

Examples

$$Y_0^0 = \frac{1}{2} \sqrt{2}, \quad Y_1^0 = \frac{1}{2} \sqrt{6} \cos \theta, \quad Y_2^0 = \frac{1}{4} \sqrt{10} (3 \cos^2 \theta - 1)$$

$$Y_1^{-1} = \frac{1}{2} \sqrt{3} \sin \theta, \quad Y_2^{-1} = \frac{1}{2} \sqrt{15} \sin \theta \cos \theta, \quad Y_2^{-2} = \frac{1}{4} \sqrt{15} \sin^2 \theta$$

Physical rationale for inequality $-m \leq m \leq \ell$



Vector of magnitude $\ell\hbar$ has a z-component given by $\ell\hbar \sin \theta$ which has a range classically of $-\ell\hbar$ to $\ell\hbar$ ($0 \leq \theta \leq \pi$).

Remarks: For each value of ℓ there are $2\ell + 1$ values of m ; there are $2\ell + 1$ spherical harmonics.

$\ell = 1$ in detail:

$$Y_{10}(\theta, \phi) = Y_1^0(\cos \theta) = \frac{1}{2} \sqrt{6} \cos \theta$$

$$Y_{1,+1}(\theta, \phi) = Y_1^{+1}(\cos \theta) \frac{e^{i\phi}}{\sqrt{2}},$$

$$\sin \theta e^{i\phi}$$

$$Y_{1,-1}(\theta, \phi) = Y_1^{-1}(\cos \theta) \frac{e^{-i\phi}}{\sqrt{2}},$$

$$\sin \theta e^{-i\phi}$$

$$L^2(Y_{1,+1} + Y_{1,-1}) = \hbar^2(Y_{1,+1} + Y_{1,-1})$$

Normalization factor

So, $\frac{(Y_{1,+1} + Y_{1,-1})}{\sqrt{2}}$ is still an eigenfunction of L_{op}^2 but *not* of L_z .

Why consider this linear combination? It's real, i.e., $\sin \theta \cos \theta$
 ($x = r \sin \theta \cos \theta$)

$$\frac{(Y_{1,+1} - Y_{1,-1})}{\sqrt{2}} \sin \theta \cos \phi \quad (y = r \sin \theta \cos \phi)$$

These orbitals are used in discussions of bonding.